

T-duality - A pedagogical introduction

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- Introduction to (topological) T-duality
- Spherical T-duality
- Non-isometric T-duality (*not covered*)

Based on collaborations with Mathai and many others (Evslin, Hannabuss, Sati, Wu, Klimčík, Bugden, Wright, ...)

Fourier Transform

Fourier series for $f : S^1 \rightarrow \mathbb{R}$

$$\widehat{f}(n) = \frac{1}{2\pi} \int_0^{2\pi} f(x) e^{-inx} dx$$
$$f(x) = \sum_{n \in \mathbb{Z}} \widehat{f}(n) e^{inx}$$

Fourier transform for $f : \mathbb{R} \rightarrow \mathbb{R}$

$$\widehat{f}(p) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) e^{-ipx} dx$$
$$f(x) = \int_{-\infty}^{\infty} \widehat{f}(p) e^{ipx} dp$$

Fourier Transform - cont'd

More generally, for G a locally compact, abelian group, we have a Fourier transform $\mathcal{F} : \text{Fun}(G) \rightarrow \text{Fun}(\widehat{G})$

$$\widehat{f}(p) = \int_G f(x) e^{-ipx} dx = \mathcal{F}(f)(p)$$
$$f(x) = \int_{\widehat{G}} \widehat{f}(p) e^{ipx} dp$$

where

$$\widehat{G} = \text{Hom}(G, U(1)) = \text{char}(G)$$

is the Pontryagin dual of G . I.e. a character is a $U(1)$ valued function on G , satisfying $\chi(x + y) = \chi(x)\chi(y)$.

The characters form a locally compact, abelian group \widehat{G} under pointwise multiplication.

$$\begin{aligned} G = S^1, & \quad \widehat{G} = \mathbb{Z}, & e^{inx} \\ G = \mathbb{R}, & \quad \widehat{G} = \mathbb{R}, & e^{ipx} \end{aligned}$$

We can think of $\chi(x, p) = e^{ipx} \in \text{Fun}(G \times \widehat{G})$ as the ‘universal’ character.

Fourier transform expresses the fact that the characters of G span $\text{Fun}(G)$.

I.e. we have the following “correspondence”

$$\begin{array}{ccc} & \mathbf{G} \times \widehat{\mathbf{G}} & \\ \pi \swarrow & & \searrow \widehat{\pi} \\ \mathbf{G} & & \widehat{\mathbf{G}} \end{array}$$

$$\mathcal{F}f = \widehat{\pi}_*(\pi^*(f) \times \chi(x, p))$$

T-duality is a geometric version of harmonic analysis, i.e. by replacing functions by geometric objects (such as bundles, sheaves, D-modules, ...) or, as an intermediate step, by topological characteristics associated to these objects (cohomology, K-theory, derived categories, ...).

Fourier-Mukai transform

Consider a manifold $P = M \times S^1$. By the Künneth theorem we have

$$H^\bullet(P) \cong H^\bullet(M) \otimes H^\bullet(S^1)$$

i.e.

$$H^n(P) \cong H^n(M) \oplus H^{n-1}(M)$$

We have a similar decomposition at the level of forms

$$\Omega^n(P)^{\text{inv}} \cong \Omega^n(M) \oplus \Omega^{n-1}(M).$$

i.e. invariant degree n forms on P are of the form ω or $\omega \wedge d\theta$, where ω is an n , respectively $n - 1$, form on M .

Consider $\widehat{P} = M \times \widehat{S}^1$. We have an isomorphism

$$\mathcal{F} : H^{\bar{i}}(P) \xrightarrow{\cong} H^{\bar{i}+1}(\widehat{P})$$

where

$$H^{\bar{0}}(P) = \bigoplus_{i \geq 0} H^{2i}(P), \quad H^{\bar{1}}(P) = \bigoplus_{i \geq 0} H^{2i+1}(P),$$

Explicitly

$$\omega \mapsto d\hat{\theta} \wedge \omega, \quad d\theta \wedge \omega \mapsto \omega$$

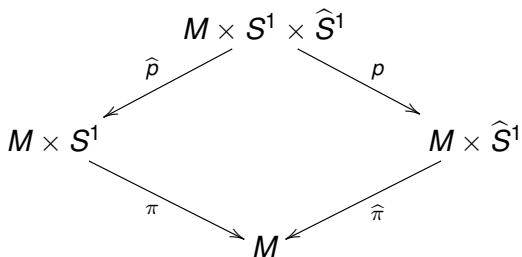
or

$$\mathcal{F}\Omega = \int_{S^1} (1 + d\theta \wedge d\hat{\theta}) \Omega = \int_{S^1} e^{d\theta \wedge d\hat{\theta}} \Omega = \int_{S^1} e^F \Omega$$

Fourier-Mukai transform - cont'd

I.e. \mathcal{F} is given by a correspondence

$$\mathcal{F}\Omega = p_* (\hat{p}^* \Omega \wedge e^F)$$



Fourier-Mukai transform - cont'd

Once we recognize that $F = d\theta \wedge d\hat{\theta}$ is the curvature of a canonical linebundle \mathcal{P} (the Poincaré linebundle) over $S^1 \times \hat{S}^1$, in fact $e^F = \text{ch}(\mathcal{P})$, this immediately suggests a 'geometrization' in terms of vector bundles over P and \hat{P} . (*)

$$\mathcal{F}E = p_* (\hat{p}^* E \otimes \mathcal{P})$$

This gives rise to the so-called Fourier-Mukai transform

$$\mathcal{F} : K^i(P) \xrightarrow{\cong} K^{i+1}(\hat{P})$$

which has many of the properties of the Fourier transform discussed earlier.

The discussion can be generalized to complexes of vector bundles (complexes of sheaves) and thus gives rise to a Fourier-Mukai correspondence between derived categories $D(P)$ and $D(\hat{P})$.

T-duality - Closed string on $M \times S^1$

Closed strings on $M \times S^1$ are described by

$$X : \Sigma \rightarrow M \times S^1$$

where $\Sigma = \{(\sigma, \tau)\}$ is the closed string worldsheet.

Upon quantization, we find

- Momentum modes: $p = \frac{n}{R}$
- Winding modes: $X(0, \tau) \sim X(1, \tau) + mR$

$$E = \left(\frac{n}{R}\right)^2 + (mR)^2 + \text{osc. modes}$$

We have a duality $R \rightarrow 1/R$, such that ST on $M \times S^1$ is equivalent to ST on $M \times \widehat{S}^1$ (or a duality between IIA and IIB ST, for susy ST)

T-duality - Principal S^1 -bundles

Suppose we have a pair (P, H) , consisting of a principal circle bundle

$$\begin{array}{ccc} S^1 & \longrightarrow & P \\ & & \downarrow \pi \\ & & M \end{array}$$

and a so-called H-flux H on P , a Čech 3-cocycle.

Topologically, P is classified by an element in $F \in H^2(M, \mathbb{Z})$ while H gives a class in $H^3(P, \mathbb{Z})$

T-duality - Principal S^1 -bundles

The (topological) T-dual of (P, H) is given by the pair $(\widehat{P}, \widehat{H})$, where the principal S^1 -bundle

$$\begin{array}{ccc} \widehat{S}^1 & \longrightarrow & \widehat{P} \\ & & \downarrow \widehat{\pi} \\ & & M \end{array}$$

and the dual H-flux $\widehat{H} \in H^3(\widehat{P}, \mathbb{Z})$, satisfy

$$\widehat{F} = \pi_* H, \quad F = \widehat{\pi}_* \widehat{H}$$

where $\pi_* : H^3(P, \mathbb{Z}) \rightarrow H^2(M, \mathbb{Z})$, is the pushforward map ('integration over the S^1 -fibre').

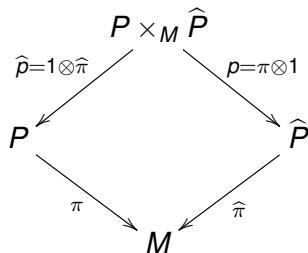
T-duality - Principal S^1 -bundles

The ambiguity in the choice of \widehat{H} is (almost) removed by requiring that

$$\widehat{p}^*H - p^*\widehat{H} \equiv 0 \in H^3(P \times_M \widehat{P}, \mathbb{Z})$$

where $P \times_M \widehat{P}$ is the correspondence space

$$P \times_M \widehat{P} = \{(x, \widehat{x}) \in P \times \widehat{P} \mid \pi(x) = \widehat{\pi}(\widehat{x})\}$$



Gysin sequences

$$\dots \longrightarrow H^3(M) \xrightarrow{\pi^*} H^3(P) \xrightarrow{\pi_*} H^2(M) \xrightarrow{\cup F} H^4(M) \longrightarrow \dots$$

$$\dots \longrightarrow H^3(M) \xrightarrow{\widehat{\pi}^*} H^3(\widehat{P}) \xrightarrow{\widehat{\pi}_*} H^2(M) \xrightarrow{\cup \widehat{F}} H^4(M) \longrightarrow \dots$$

T-duality - Principal S^1 -bundles

$$\begin{array}{cccccccc}
 0 & \xrightarrow{\cup \widehat{F}} & H^1(M) & \xrightarrow{\widehat{\pi}^*} & H^1(\widehat{P}) & \xrightarrow{\widehat{\pi}_*} & H^0(M) & \xrightarrow{\cup \widehat{F}} & H^2(M) & \longrightarrow & \dots \\
 \downarrow \cup F & & \downarrow \cup F & & \downarrow \cup \widehat{\pi}^* F & & \downarrow \cup F & & \downarrow \cup F & & \\
 H^1(M) & \xrightarrow{\cup \widehat{F}} & H^3(M) & \xrightarrow{\widehat{\pi}^*} & H^3(\widehat{P}) & \xrightarrow{\widehat{\pi}_*} & H^2(M) & \xrightarrow{\cup \widehat{F}} & H^4(M) & \longrightarrow & \dots \\
 \downarrow \pi^* & & \downarrow \pi^* & & \downarrow p^* & & \downarrow \pi^* & & \downarrow \pi^* & & \\
 H^1(P) & \xrightarrow{\cup \pi^* \widehat{F}} & H^3(P) & \xrightarrow{\widehat{p}^*} & H^3(P \times_M \widehat{P}) & \xrightarrow{\widehat{p}_*} & H^2(P) & \xrightarrow{\cup \pi^* \widehat{F}} & H^4(P) & \longrightarrow & \dots \\
 \downarrow \pi_* & & \downarrow \pi_* & & \downarrow p_* & & \downarrow \pi_* & & \downarrow \pi_* & & \\
 H^0(M) & \xrightarrow{\cup \widehat{F}} & H^2(M) & \xrightarrow{\widehat{\pi}^*} & H^2(\widehat{P}) & \xrightarrow{\widehat{\pi}_*} & H^1(M) & \xrightarrow{\cup \widehat{F}} & H^3(M) & \longrightarrow & \dots \\
 \downarrow & & \downarrow & & \downarrow & & \downarrow & & \downarrow & & \\
 \vdots & & \vdots & & \vdots & & \vdots & & \vdots & &
 \end{array}$$

Consider principal S^1 -bundles P over $M = S^2$, then

$$H^2(M, \mathbb{Z}) \cong \mathbb{Z}, \quad H^3(P, \mathbb{Z}) \cong \mathbb{Z}$$

and we have, for example,

$$(S^2 \times S^1, 0) \longrightarrow (S^2 \times S^1, 0)$$

$$(S^2 \times S^1, 1) \longrightarrow (S^3, 0)$$

or more generally

$$(L_p, k) \longrightarrow (L_k, p)$$

where $L_p = S^3/\mathbb{Z}_p$ is the lens space.

T-duality - Twisted cohomology

Using $\Omega^k(P)^{inv} \cong \Omega^k(M) \oplus \Omega^{k-1}(M)$

$$F = dA, \quad H = H_{(3)} + A \wedge H_{(2)}$$

we find

$$\widehat{F} = H_{(2)} = d\widehat{A}, \quad \widehat{H} = H_{(3)} + \widehat{A} \wedge F$$

such that

$$\widehat{H} - H = \widehat{A} \wedge F - A \wedge \widehat{F} = d(A \wedge \widehat{A}).$$

Theorem

We have an isomorphism of (\mathbb{Z}_2 -graded) differential complexes

$$T_* : (\Omega(P)^{inv}, d_H) \longrightarrow (\Omega(\widehat{P})^{inv}, d_{\widehat{H}})$$

where $d_H = d + H \wedge$.

Proof.

Define

$$T_*\omega = \int_{S^1} e^{A \wedge \hat{A}} \omega$$

then

$$d_H T_* = T_* d_{\hat{H}}.$$



and consequently, we have isomorphisms

$$T_* : H^{\bar{i}}(P, H) \xrightarrow{\cong} H^{\bar{i}+1}(\hat{P}, \hat{H})$$

as well as

$$T_* : K^i(P, H) \xrightarrow{\cong} K^{i+1}(\widehat{P}, \widehat{H})$$

For example,

$$K^i(L_p, k) \cong \begin{cases} \mathbb{Z}_k & i = 0 \\ \mathbb{Z}_p & i = 1 \end{cases}$$

The physics

	String Theory $M_4 \times Y_6$	
$\mathcal{N} = 1$ $\mathcal{N} = 2$ $\mathcal{N} = 3$	Complex manifold Kähler Calabi-Yau Hyper-Kähler	
	S^1 Strings $H \in H^3(Y, \mathbb{Z})$ Mirror Symmetry / T-duality generalized geometry	
	$S^1 \longrightarrow S^3$ \downarrow S^2	

The physics

	String Theory $M_4 \times Y_6$	M-Theory / 11D SUGRA $M_4 \times Y_7$
$\mathcal{N} = 1$ $\mathcal{N} = 2$ $\mathcal{N} = 3$	Complex manifold Kähler Calabi-Yau Hyper-Kähler	Contact manifold Sasakian Sasaki-Einstein 3-Sasakian
	S^1 Strings $H \in H^3(Y, \mathbb{Z})$ Mirror Symmetry / T-duality generalized geometry	S^3 2- and 5-branes $H \in H^7(Y, \mathbb{Z})$ Spherical T-duality? M-geometry?
	$S^1 \longrightarrow S^3$ \downarrow S^2	$S^3 \longrightarrow S^7$ \downarrow S^4

Spherical T-duality - Principal SU(2)-bundles

Much of the above can be generalized to principal SU(2)-bundles:

Gysin sequence for principal SU(2)-bundles $\pi : P \rightarrow M$

$$\dots \longrightarrow H^7(M) \xrightarrow{\pi^*} H^7(P) \xrightarrow{\pi_*} H^4(M) \xrightarrow{\cup c_2(P)} H^8(M) \longrightarrow \dots$$

where

$$c_2(P) = \frac{1}{8\pi^2} \text{Tr}(F \wedge F) \in H^4(M)$$

is (a de Rham representative of) the 2nd Chern class of P . However, in this case,

$$[M, BSU(2)] \longrightarrow H^4(M, \mathbb{Z})$$

is, in general, neither surjective nor injective.

Physical derivation of T-duality

A 2D non-linear sigma model describes maps X from a 2-dimensional surface ('worldsheet') Σ to an N -dimensional manifold M ('target'), equipped with additional structure

For example

$$S[X] = \frac{1}{2} \int_{\Sigma} G_{ij}(X) dX^i \wedge \star dX^j + B_{ij}(X) dX^i \wedge dX^j$$

Symmetries of sigma model

Given a set of vector fields $v_a(X) = v_a^i(X)\partial_i$ forming a Lie algebra \mathfrak{g}

$$[v_a, v_b] = C^c{}_{ab}v_c$$

Consider the infinitesimal transformations

$$\delta_\epsilon X^i = v_a^i(X) \epsilon^a$$

we have

$$\delta_\epsilon S = \int_\Sigma \epsilon^a \left((\mathcal{L}_{v_a} G)_{ij} dX^i \wedge \star dX^j + (\mathcal{L}_{v_a} B)_{ij} dX^i \wedge dX^j \right)$$

The sigma model action is invariant under these transformations if

$$\mathcal{L}_{v_a} G = 0, \quad \mathcal{L}_{v_a} B = 0$$

If this is the case, we can *gauge* the model by promoting the global symmetry to a local one (i.e. take $\epsilon \in C^\infty(\Sigma, \mathfrak{g})$)

The gauged action

Introducing gauge fields $A \in \Omega^1(\Sigma, \mathfrak{g})$ the gauged action is given by

$$S[X, A] = \frac{1}{2} \int_{\Sigma} G_{ij}(X) DX^i \wedge \star DX^j + B_{ij}(X) DX^i \wedge DX^j$$

where

$$DX^i = dX^i - v_a^i A^a$$

are the covariant derivatives.

Gauge invariance

The gauged action $S[X, A]$ is invariant with respect to the following (local) gauge transformations:

$$\delta_\epsilon X^i = v_a^i \epsilon^a$$

$$\delta_\epsilon A = d\epsilon + [A, \epsilon] = (d\epsilon^a + C^a_{bc} A^b \epsilon^c) T_a$$

where T_a is a basis of \mathfrak{g} .

Now suppose we want the gauged sigma model to be equivalent to the the ungauged model. Then we need to 'fix the gauge'

Introduce the curvature $F \in \Omega^2(\Sigma, \mathfrak{g})$

$$F = dA + A \wedge A = (dA^a + \frac{1}{2}C^a_{bc} A^b \wedge A^c) T_a = F^a T_a$$

and an 'auxiliary field' $\widehat{X} \in C^\infty(\Sigma, \mathfrak{g}^*)$, with infinitesimal transformation rules

$$\begin{aligned}\delta_\epsilon F^a &= C^a_{bc} F^b \epsilon^c \\ \delta_\epsilon \widehat{X}_a &= -C^c_{ab} \widehat{X}_c \epsilon^b\end{aligned}$$

Consider the action

$$S[X, A, \hat{X}] = \frac{1}{2} \int_{\Sigma} (G_{ij}(X) DX^i \wedge \star DX^j + B_{ij}(X) DX^i \wedge DX^j) \\ + \int_{\Sigma} \hat{X}_a F^a$$

The equation of motion for \hat{X}_a gives $F^a = 0$.

To solve this equation we need to lift the action of \mathfrak{g} to an action of the group G ($\mathfrak{g} = \text{Lie } G$)

Example: Group manifold

Let $g : \Sigma \rightarrow G$

$$S[g] = \frac{1}{2} \int_{\Sigma} (g^{-1} dg \wedge *g^{-1} dg)_G$$

Invariant under left action of $h \in G$

$$S[hg] = S[g]$$

while

$$S[gh] = \frac{1}{2} \int_{\Sigma} (\text{Ad}(h^{-1})g^{-1} dg \wedge \text{Ad}(h^{-1}) * g^{-1} dg)_G$$

So, invariant under right action of G if G is Ad-invariant (Killing form)

Example: Gauged

In that case we can gauge in the standard way, and obtain the gauged model (with F -term)

$$S[g, A, \hat{X}] = \frac{1}{2} \int_{\Sigma} (g^{-1} Dg \wedge *g^{-1} Dg)_G + \int_{\Sigma} \langle \hat{X}, F \rangle$$

where

$$\begin{aligned} g^{-1} Dg &= g^{-1} dg - A \\ F &= dA + A \wedge A \end{aligned}$$

and gauge symmetry, for $h \in G$

$$\begin{aligned} g &\rightarrow gh \\ A &\rightarrow h^{-1} Ah + h^{-1} dh \\ \hat{X} &\rightarrow \text{Ad}^*(h^{-1})\hat{X} \end{aligned}$$

Example: Gauged model

Solving $F = 0$ gives $A = -dkk^{-1}$ for $k \in C^\infty(\Sigma, G)$, and substituting

$$g^{-1}Dg \rightarrow g^{-1}dg + dkk^{-1} = k((gk)^{-1}d(gk))k^{-1}$$

i.e.

$$S[g, A = -dkk^{-1}] = S[gk]$$

so after 'fixing the gauge' we recover the ungauged model.

On the other hand, first solving the equation of motion for A , and then fixing the gauge, gives dual model

$$\widehat{S}[\widehat{X}] = \frac{1}{2} \int_{\Sigma} \widehat{G}^{ab}(\widehat{X}) d\widehat{X}_a \wedge \star d\widehat{X}_b$$

with dual 'metric'

$$\widehat{G}^{-1}_{ab} = G_{ab} - C^c_{ab} \widehat{X}_c$$

Suppose we have a $U(1)^N$ isometry $X^m \rightarrow X^m + \epsilon^m$, then the T-duality rules are given by the **Buscher rules**

$$\begin{aligned}\widehat{Q}_{ij} &= \begin{pmatrix} \widehat{Q}_{\mu\nu} & \widehat{Q}_{\mu n} \\ \widehat{Q}_{m\nu} & \widehat{Q}_{mn} \end{pmatrix} \\ &= \begin{pmatrix} Q_{\mu\nu} - Q_{\mu m} (Q^{-1})^{mn} Q_{n\nu} & -Q_{\mu m} (Q^{-1})^m{}_n \\ (Q^{-1})_m{}^n Q_{n\nu} & (Q^{-1})_{mn} \end{pmatrix}\end{aligned}$$

where $Q_{ij} = G_{ij} + B_{ij}$.

More explicitly, for a $U(1)$ isometry,

$$\hat{G}_{\bullet\bullet} = \frac{1}{G_{\bullet\bullet}}$$

$$\hat{G}_{\bullet\mu} = \frac{B_{\bullet\mu}}{G_{\bullet\bullet}}$$

$$\hat{G}_{\mu\nu} = G_{\mu\nu} - \frac{1}{G_{\bullet\bullet}} (G_{\bullet\mu} G_{\bullet\nu} - B_{\bullet\mu} B_{\bullet\nu})$$

$$\hat{B}_{\bullet\mu} = \frac{G_{\bullet\mu}}{G_{\bullet\bullet}}$$

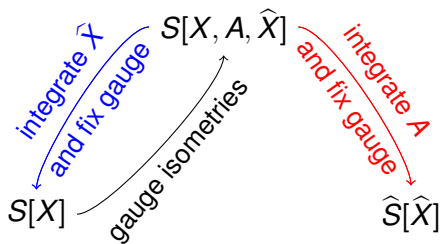
$$\hat{B}_{\mu\nu} = B_{\mu\nu} - \frac{1}{G_{\bullet\bullet}} (G_{\bullet\mu} B_{\bullet\nu} - G_{\bullet\nu} B_{\bullet\mu})$$

Letting

$$G = G_{\mu\nu} dX^\mu \otimes dX^\nu + (d\theta + A_\mu dX^\mu) \otimes (d\theta + A_\nu dX^\nu)$$

$$B = \frac{1}{2} B_{\mu\nu} dX^\mu \wedge dX^\nu + B_{\mu\bullet} dX^\mu \wedge (d\theta + A_\nu dX^\nu)$$

gives back the T-duality rules for $H = dB$ and $F = dA$ as discussed before.



THANKS